Topographical fingerprints of many-body interference blocking in STM junctions on thin insulating films

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STM on thin insulating films



Weak tip-molecule tunnelling coupling Low molecule-substrate hybridization

sequential tunnelling



URVisualization of molecular orbitals

Topography

Spectroscopy



J. Repp and G. Meyer, Physical Review Letters 94, 026803 (2005)



TR Intramolecular interference: theoretical proposals



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R. Stadler, et al. Nanotechnology, **14**, 138 (2003)



D. V. Cardamone, et al. Nano Lett., **6**, 2422 (2006)



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S.H. Ke, et al. Nano Lett., **8**, 3257 (2008)



T. Markussen, et al. Nano Lett., **10**, 4260 (2010)





Experimental evidence



Guédon et al. Nature Nanotech. 7, 305 (2012)

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Interference and dephasing

TR





Resonance

Constant height current maps in resonance and interference blockade

-0.5

0

 $V_{h}[V]$

0.5

Donarini, Siegert, Sobczyk and Grifoni Phys. Rev. B 86, 155451 (2012)

-50

Interference

blockade





Interference fingerprints



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Tunnelling dynamics



Sobczyk, Donarini, Grifoni Phys. Rev. B 85, 205408 (2012)





Many-body rate matrix

The current is proportional to the transition rate between many-body states

$$R_{N E_{0} \to N+1 E_{1}}^{\chi\tau} = \sum_{ij} N+1E_{1} |d_{i\tau}^{\dagger}|NE_{0}\rangle \Gamma_{ij}^{\chi}(E_{1}-E_{0}) \times \langle NE_{0}|d_{j\tau}N+1E_{1}\rangle f^{+}(E_{1}-E_{0}-\mu_{\chi})$$

where

$$\Gamma_{ij}^{\chi}(E_1 - E_0) = \frac{2\pi}{\hbar} \sum_k (t_{ki}^{\chi})^* t_{kj}^{\chi} \delta(\epsilon_k^{\chi} - E_1 + E_0)$$

For uncorrelated and non-degenerate systems the many-body rate reduces to

 $\epsilon_{\rm orb}$

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$$R_{N E_0 \to N+1 E_1}^{\chi \tau} = \Gamma_{\rm orb}^{\chi}(\epsilon_{\rm orb}) f^+(\epsilon_{\rm orb} - \mu_{\chi})$$

The constant current map is the isosurface of a specific molecular orbital.



Dynamics in energy space



Interference: decoupling basis

Degenerate anionic ground state

 $\ell = +1$

 $\ell = -1$

Matrix form for the many-body tunnelling rate between the neutral and anionic ground states.

Angular momentum basis

Цр

Substrate

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$$\mathbf{R}^T = R_0^T \left(\begin{array}{c} 1 \\ \mathrm{e}^{+2\mathrm{i}\phi} \end{array} \right) \stackrel{\mathrm{e}^{-2\mathrm{i}\phi}}{1}$$

Mixes angular momentum

$$\mathbf{R}^{S} = R_{0}^{S} \left(\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array} \right)$$

Conserves angular momentum

Decoupling basis

$$\tilde{\mathbf{R}}^T = R_0^T \left(\begin{array}{cc} 2 & 0\\ 0 & 0 \end{array} \right)$$

One of the anionic state is **decoupled** from the tip

2013

 $\tilde{\mathbf{R}}^S = R_0^S \left(\begin{array}{cc} 1 & 0\\ 0 & 1 \end{array} \right)$

Notice that the decoupling basis depends on the tip position.

TR Interference: current blocking





A new bottle-neck process



The **depopulation** of the blocking state via a **substrate transition** dominates the transport.





Conclusions

- We developed a semi-quantitative model for the description of double barrier STM junctions with πconjugated molecules.
- The dynamics is described in terms of many-body transitions.
- Transport through degenerate states is associated to electron interference blockade at negative sample biases.
- Close to the interference blocking regime, substrate tunnelling dominates the transport and gives flat constant height current maps.







Interference Resonance







Donarini, Begemann, and Grifoni, Phys. Rev. B 82, 125451 (2010)





Thanks



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Thank you for your attention...

