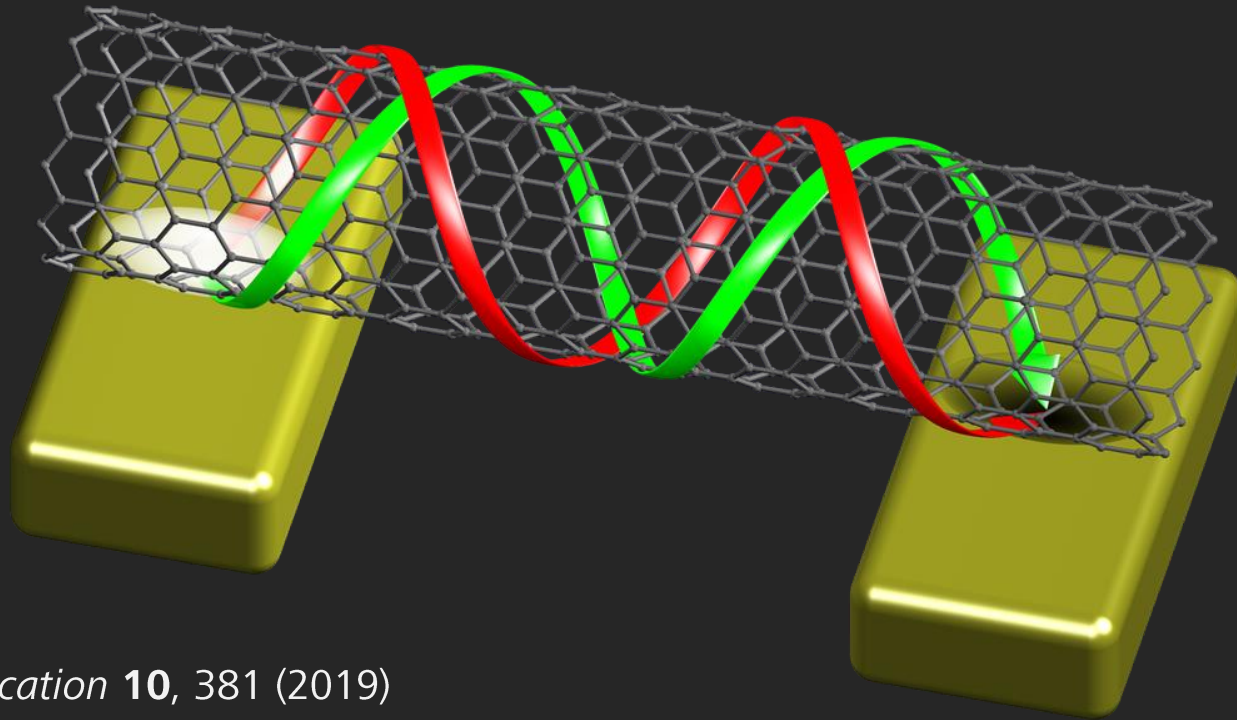


# Dark states in a carbon nanotube quantum dot



*Nature Communication* **10**, 381 (2019)



Andrea Donarini

Telluride  
30.07.2019

University of Regensburg, Germany

G. Alzetta et al. *Nuovo Cimento*. **36**, 5 (1976),

E. Arimondo and G. Orriols, *Lettere al Nuovo Cimento*, **17**, 33, (1976)

**Nonabsorbing Atomic Coherences by Coherent Two-Photon Transitions in a Three-Level Optical Pumping.**

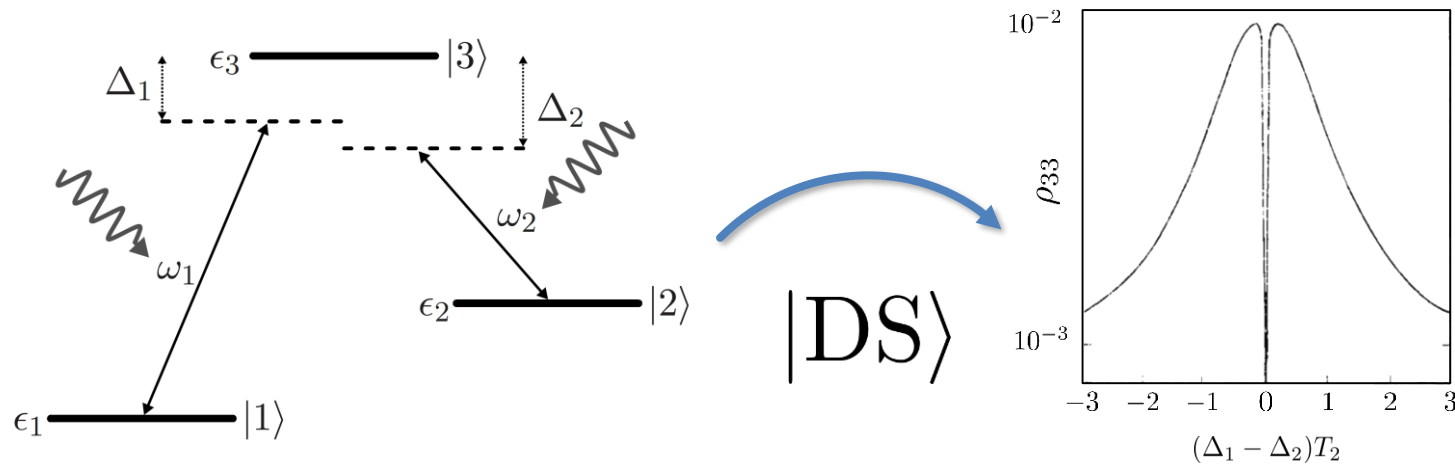
E. ARIMONDO

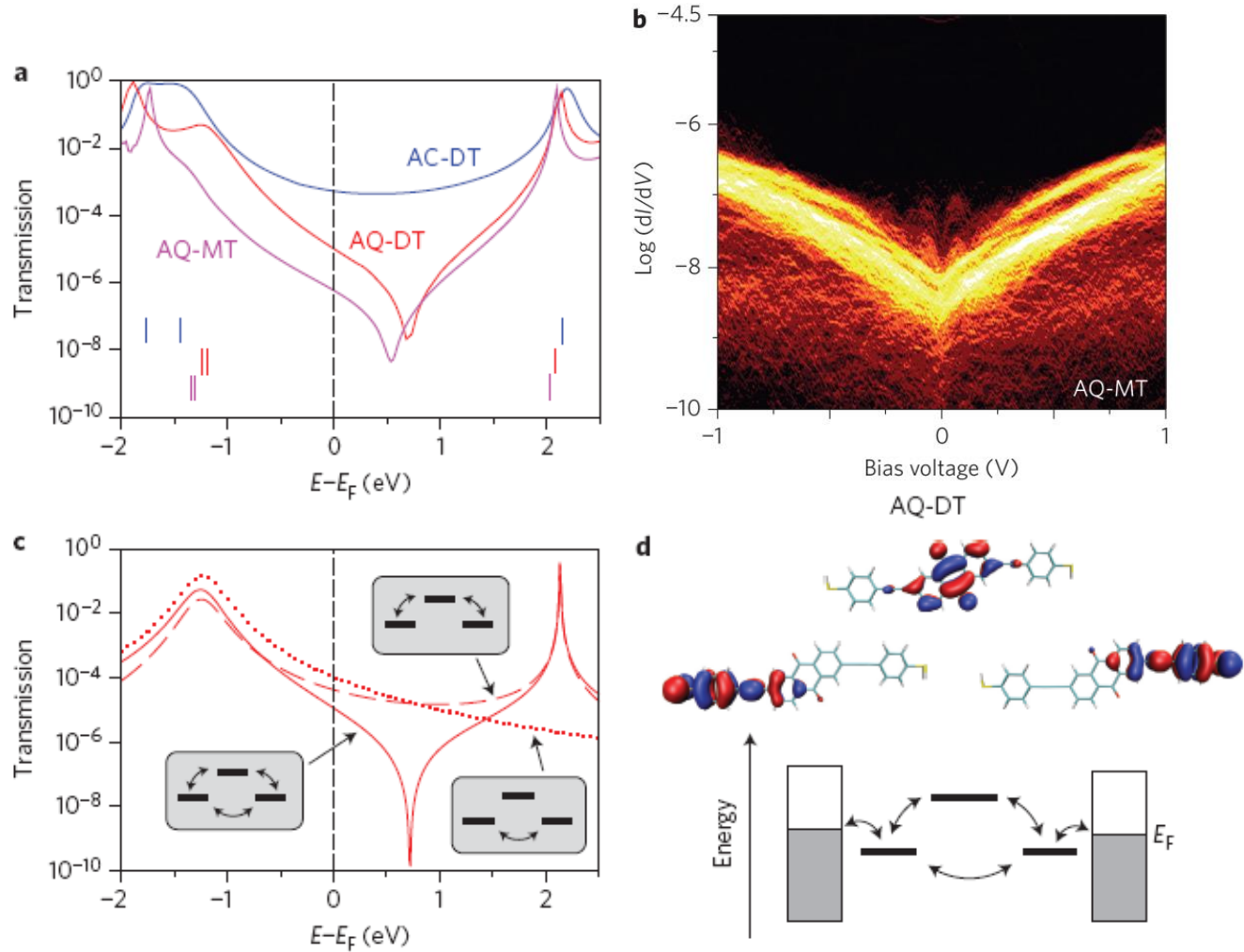
*Istituto di Fisica dell'Università - Pisa*

G. ORRIOLS

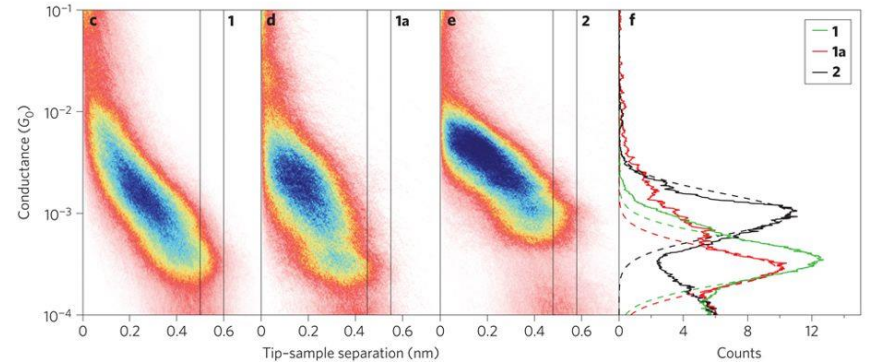
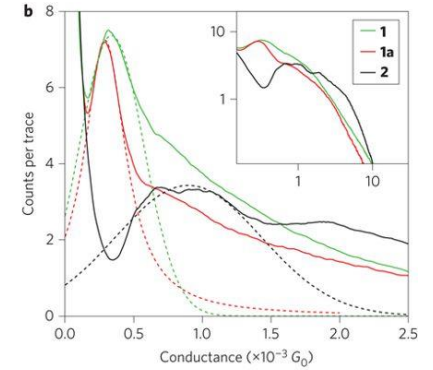
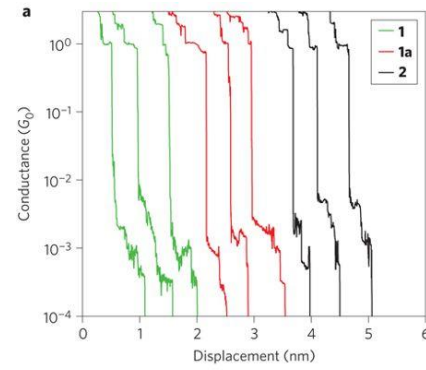
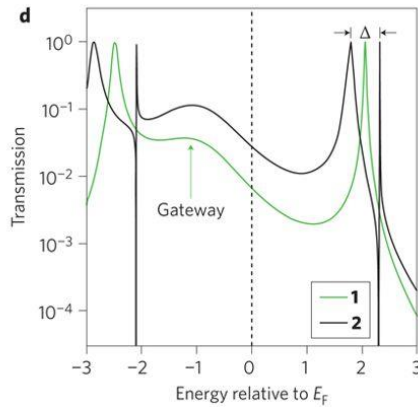
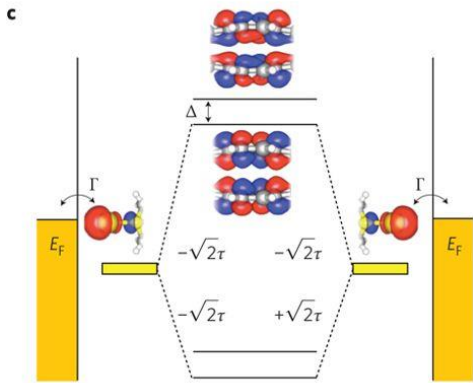
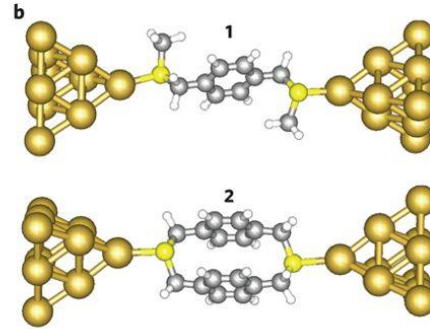
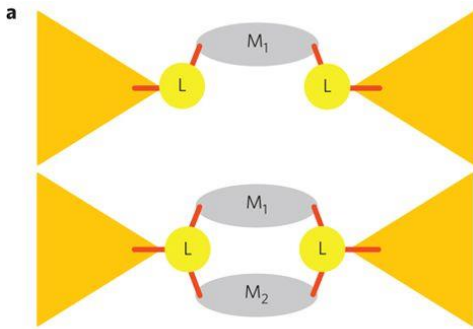
*Laboratorio di Fisica Atomica e Molecolare del C.N.R. - Pisa*

*Departament d'Òptica, Universitat de Barcelona - Barcelona*





- Zero bias conductance
- Strong coupling to the leads  $\rightarrow$  coherent (single-particle) transport

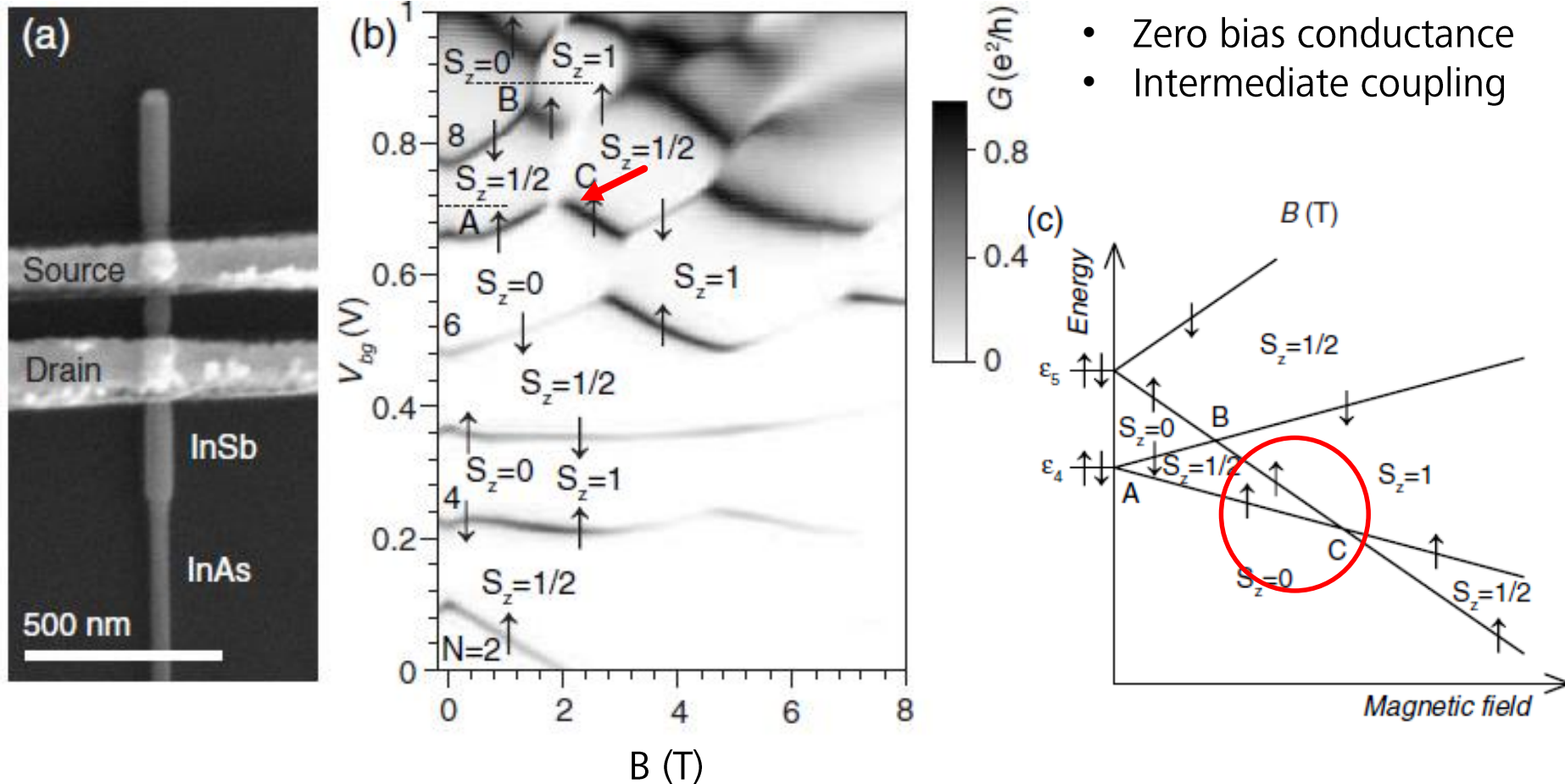


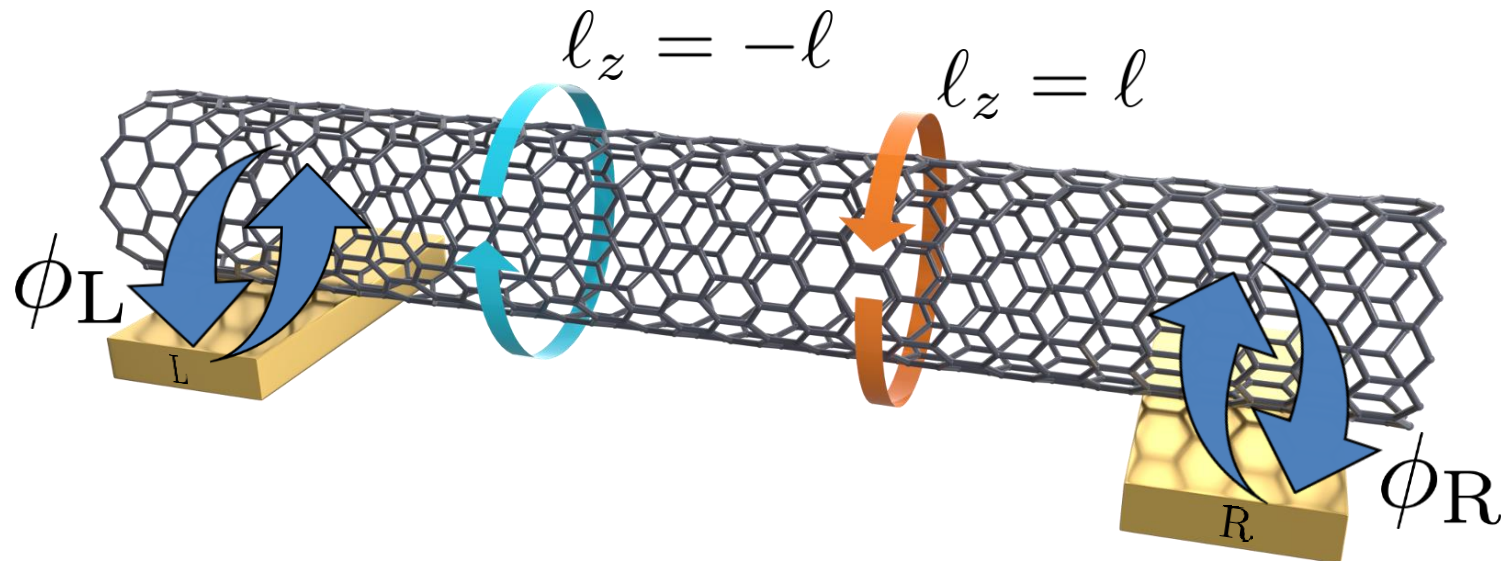
# Correlation-Induced Conductance Suppression at Level Degeneracy in a Quantum Dot

H. A. Nilsson, O. Karlström, M. Larsson, P. Caroff, J. N. Pedersen, L. Samuelson, A. Wacker,  
L.-E. Wernersson, and H. Q. Xu\*

*Nanometer Structure Consortium, Lund University, Box 118, 221 00 Lund, Sweden*

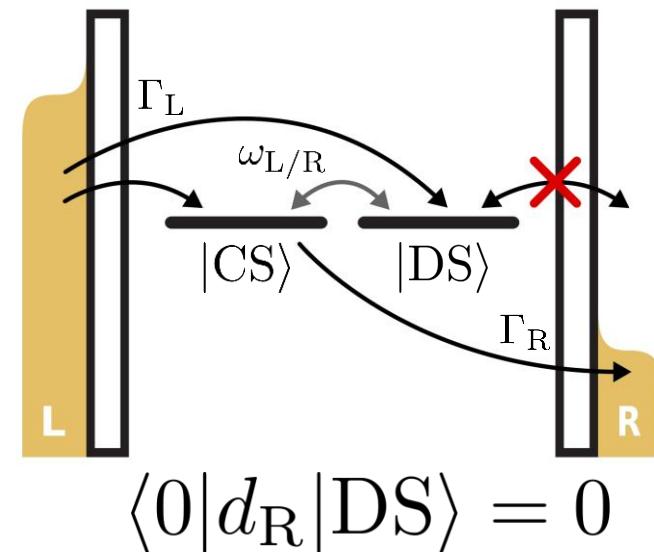
(Received 10 November 2009; published 4 May 2010)

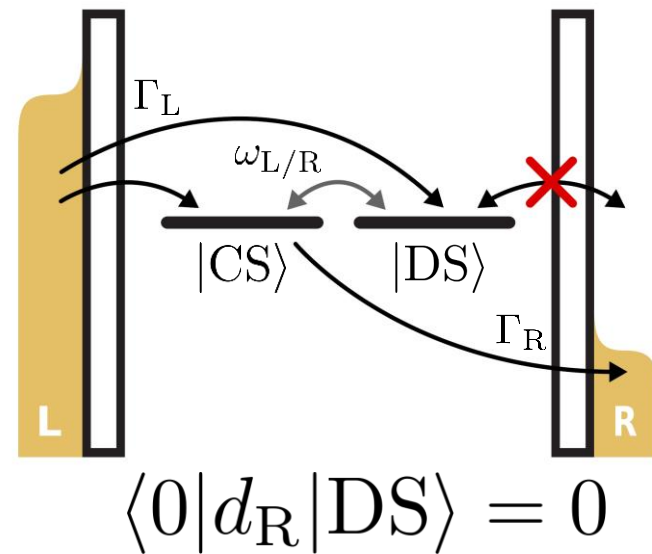
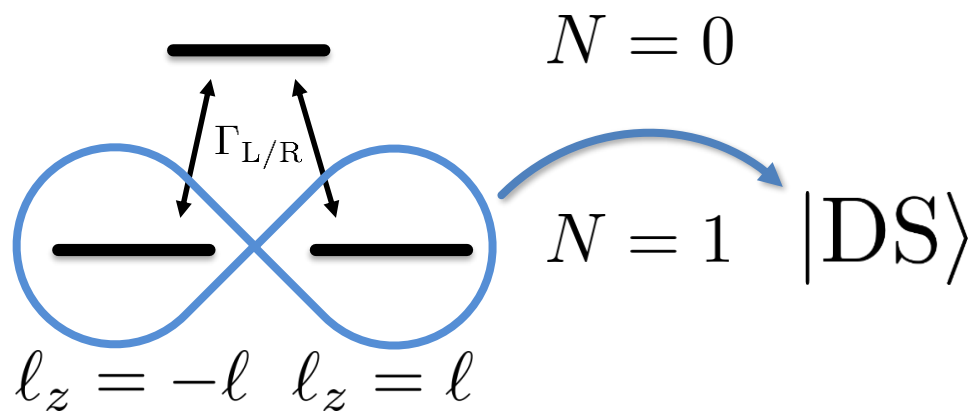
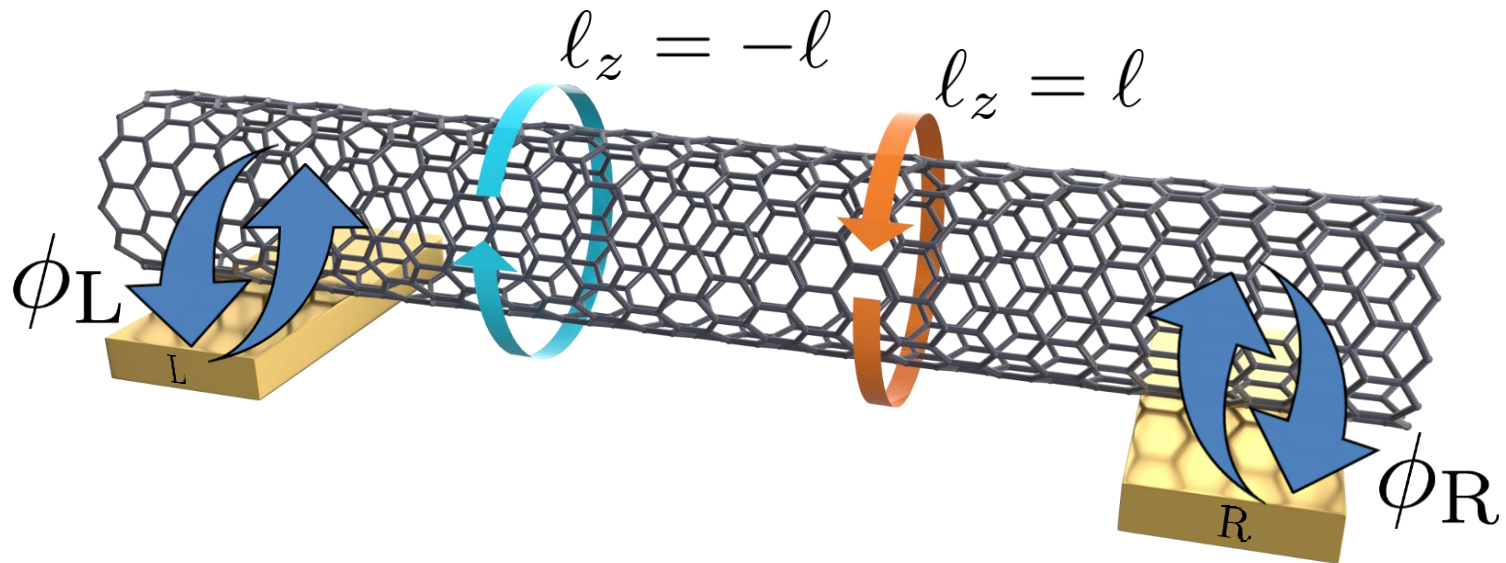




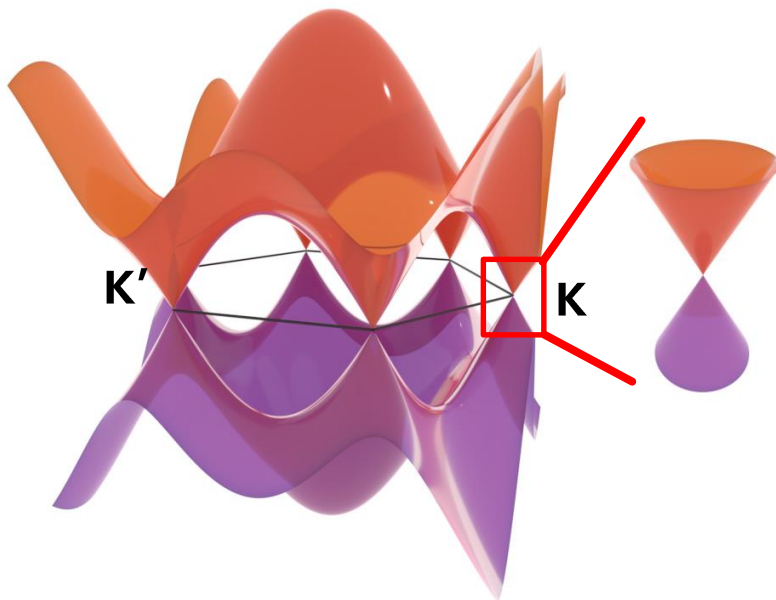
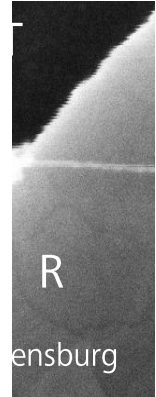
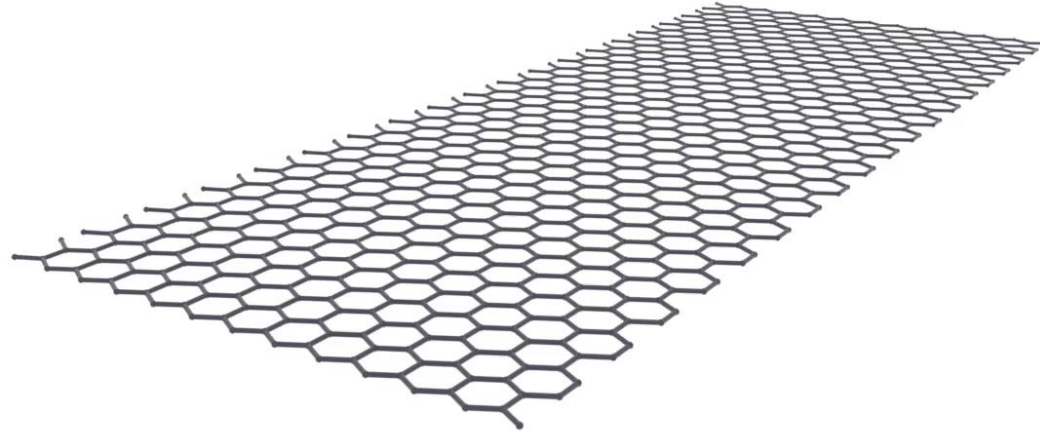
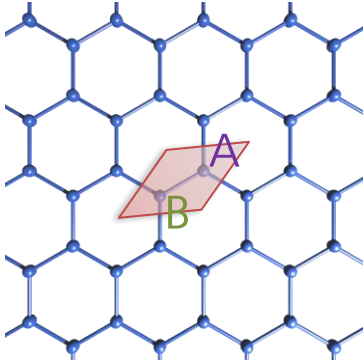
Ingredients:

- Orbital (quasi)-degeneracy
- Weak coupling to the leads  $\rightarrow$  charging effects
- Phase coherent tunnelling dynamics
- Finite bias  $\rightarrow$  directed transport





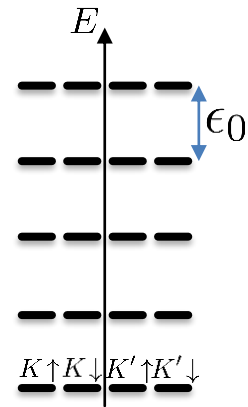
Graphene



Transverse + Longitudinal quantization

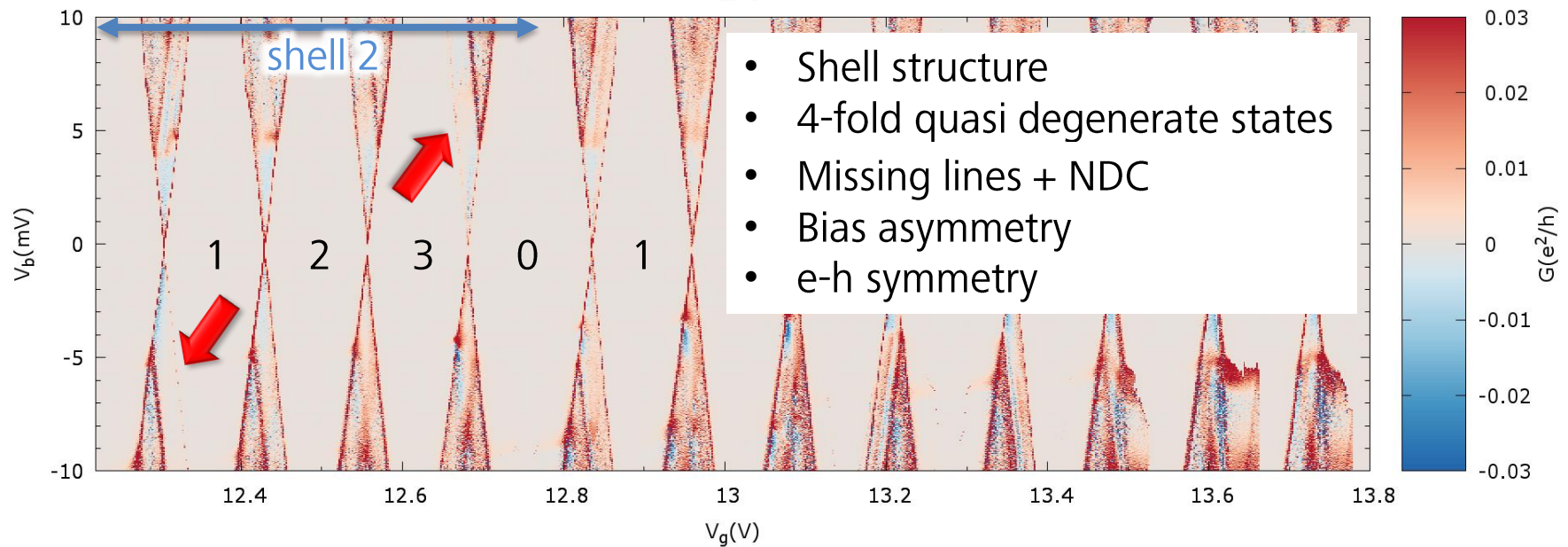
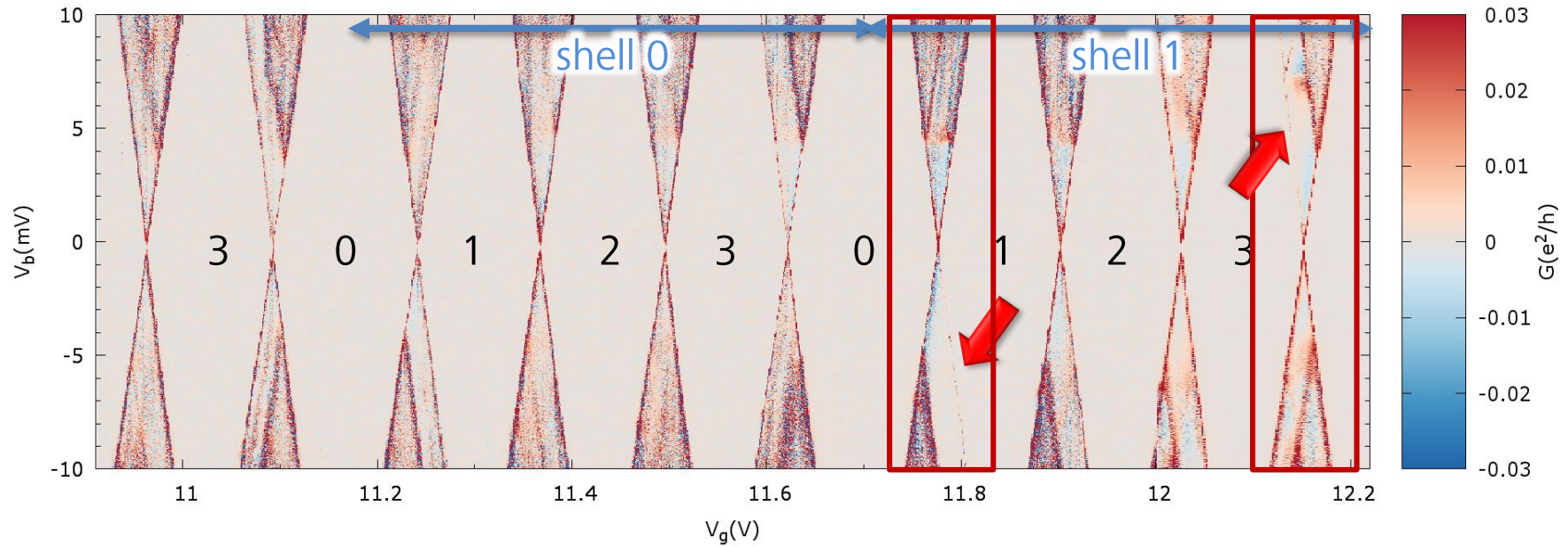


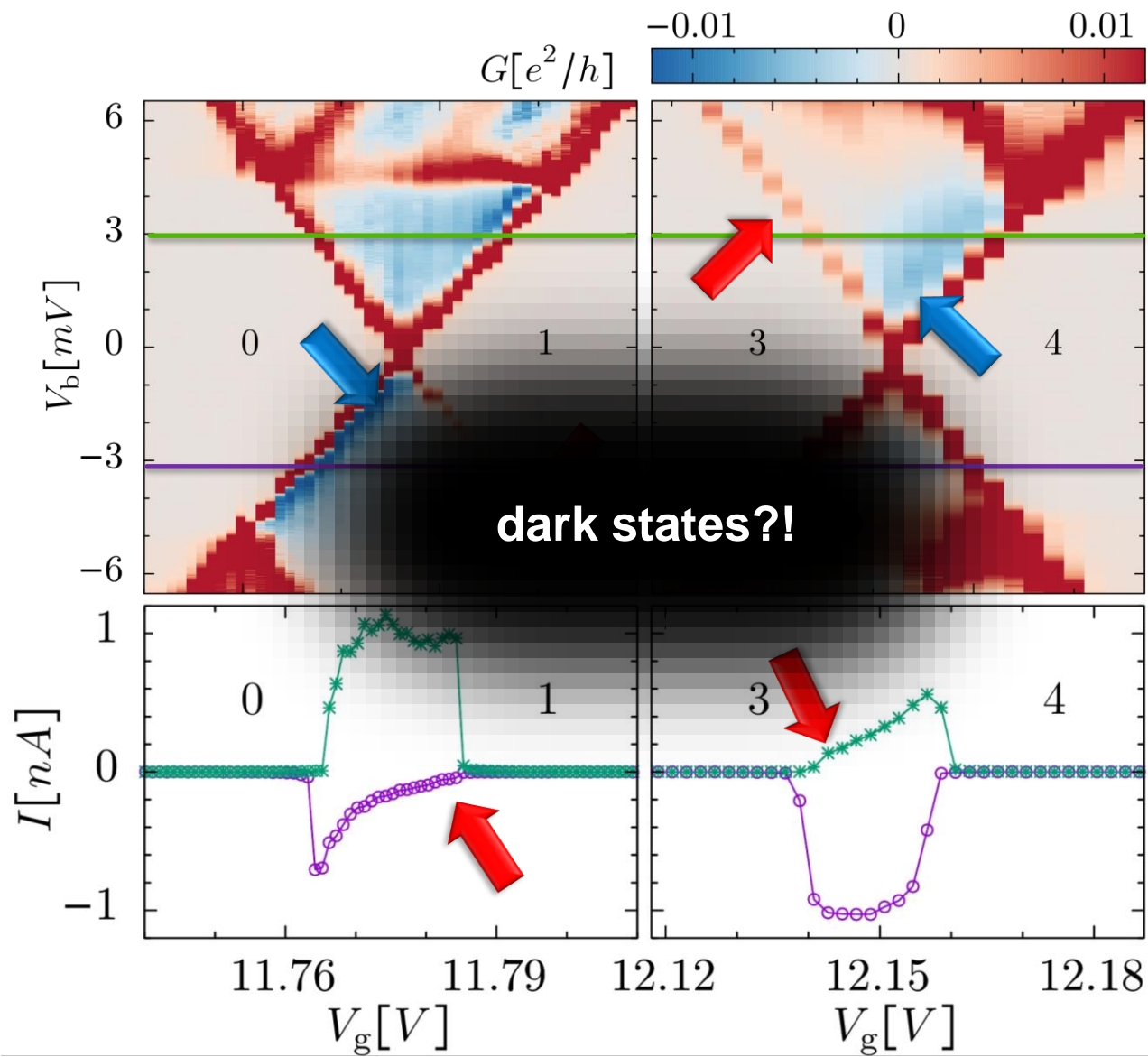
Quantum dot:

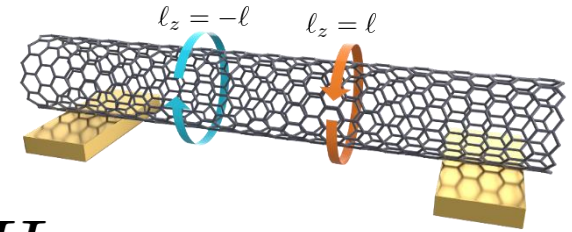


- Shell structure
- 4-fold degenerate:
  - 2x spin
  - 2x angular momentum (valley)









$$H = H_{\text{CNT}} + H_{\text{leads}} + H_{\text{tun}}$$

$$H_{\text{CNT}} = \sum_{ml_z} (m\epsilon_0 - \xi) \hat{n}_{ml_z} + \frac{U}{2} \hat{N}^2 + J \sum_m \left( \hat{\mathbf{S}}_{ml} \cdot \hat{\mathbf{S}}_{m-l} + \frac{1}{4} \hat{n}_{ml} \hat{n}_{m-l} \right)$$

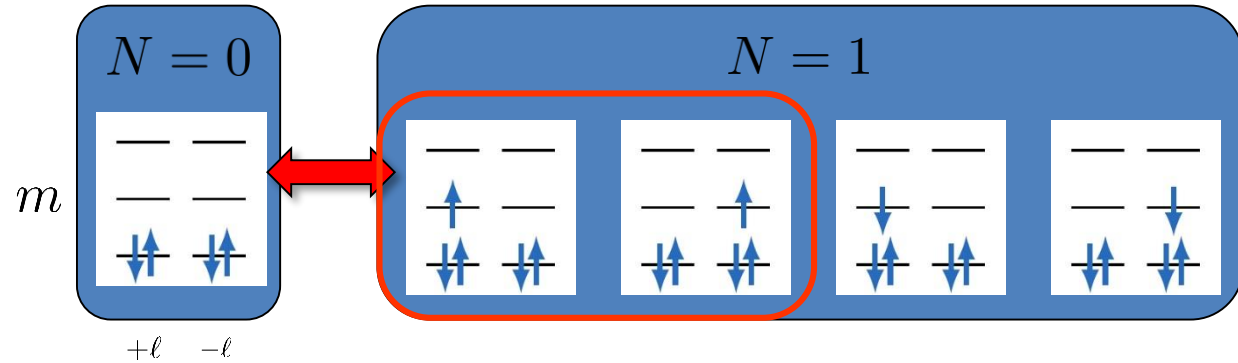
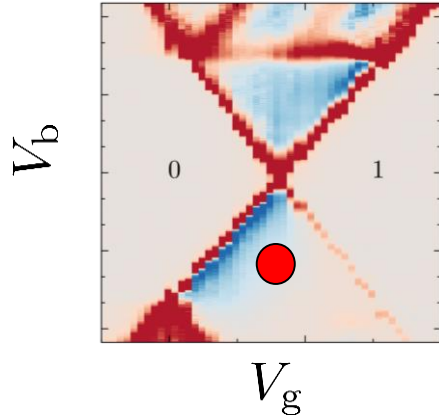
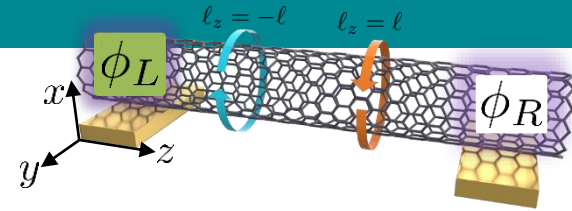
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Level spacing     $V_g$     Constant interaction    Exchange interaction for zig-zag class CNTs

$$H_{\text{leads}} = \sum_{\alpha\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\alpha\mathbf{k}\sigma}^\dagger c_{\alpha\mathbf{k}\sigma}$$

$$H_{\text{tun}} = \sum_{\alpha\mathbf{k}ml_z\sigma} t_{\alpha\mathbf{k}ml_z} d_{ml_z\sigma}^\dagger c_{\alpha\mathbf{k}\sigma} + \text{h.c.}$$

Includes the geometry of the contacts. Is treated perturbatively



$$\Gamma_{l_z l'_z}^\alpha(E_1 - E_0) = \frac{2\pi}{\hbar} \sum_{\mathbf{k}} (t_{\alpha \mathbf{k} l_z})^* t_{\alpha \mathbf{k} l'_z} \delta(\epsilon_{\mathbf{k}} - E_1 - E_0)$$



$$\Gamma^\alpha = \Gamma^\alpha \underbrace{\begin{pmatrix} 1 & ae^{2il\phi_\alpha} \\ ae^{-2il\phi_\alpha} & 1 \end{pmatrix}}_{\mathcal{R}_l : \text{coherence matrix}}$$

Atomically localized tunneling  
or  
Surface  $\Gamma$ -point approximation  
(i.e.  $k_x \approx k_F$   $k_y, k_z = 0$ )

$$a = 1$$

$$\mathcal{R}_R = \begin{pmatrix} 1 & e^{i2\phi_R} \\ e^{-i2\phi_R} & 1 \end{pmatrix} \xrightarrow{\text{diagonalize}} \mathcal{R}_R = \begin{pmatrix} 0 & 0 \\ 0 & 2 \end{pmatrix}$$

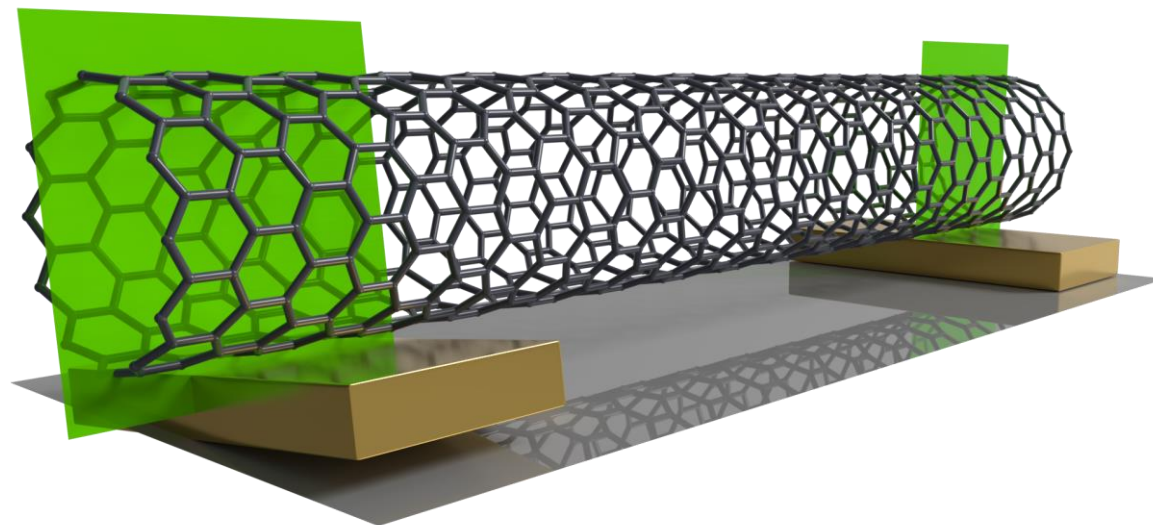
$$\mathcal{R}_L = \begin{pmatrix} 1 & e^{i2\phi_L} \\ e^{-i2\phi_L} & 1 \end{pmatrix} \xrightarrow{\Delta\phi = \phi_L - \phi_R} \mathcal{R}_L = \begin{pmatrix} 1 - \cos \Delta\phi & -i \sin \Delta\phi \\ i \sin \Delta\phi & 1 + \cos \Delta\phi \end{pmatrix}$$

Cannot be diagonalized at the same time if  $\Delta\phi \neq n\pi$

The dark state is defined with respect to a **specific** lead

$$|DS, \uparrow \alpha\rangle = \frac{1}{\sqrt{2}} \left( e^{i l \phi_\alpha} \begin{array}{c} \text{---} \text{---} \\ \uparrow \text{---} \\ \uparrow\uparrow \uparrow\uparrow \end{array} - e^{-i l \phi_\alpha} \begin{array}{c} \text{---} \text{---} \\ \text{---} \uparrow \\ \uparrow\uparrow \uparrow\uparrow \end{array} \right)$$

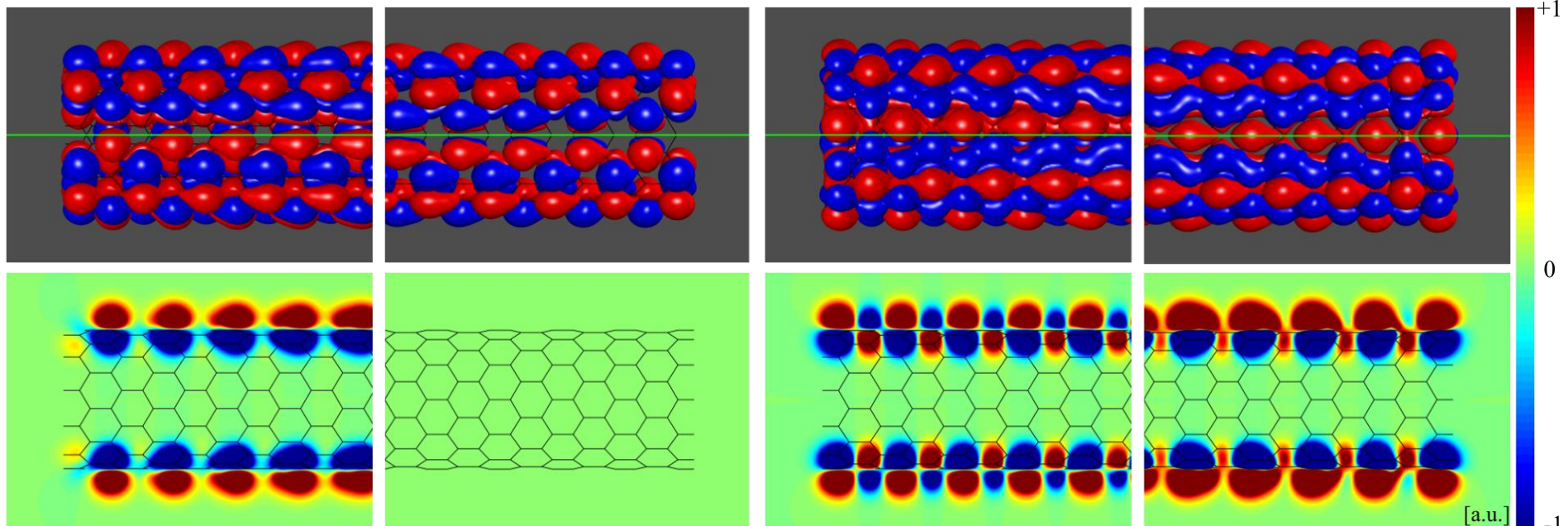
$$|CS, \uparrow \alpha\rangle = \frac{1}{\sqrt{2}} \left( e^{i l \phi_\alpha} \begin{array}{c} \text{---} \text{---} \\ \uparrow \text{---} \\ \uparrow\uparrow \uparrow\uparrow \end{array} + e^{-i l \phi_\alpha} \begin{array}{c} \text{---} \text{---} \\ \text{---} \uparrow \\ \uparrow\uparrow \uparrow\uparrow \end{array} \right)$$



(12,0) CNT

Dark state

Coupled state



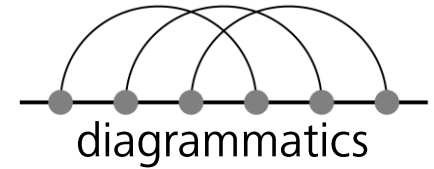
Weak coupling

Increase coupling

Strong coupling

Density matrix approach

$$\mathcal{O}(\Gamma) \dots \mathcal{O}(\Gamma^2) \dots$$



Current

$$\langle \dot{N} \rangle$$

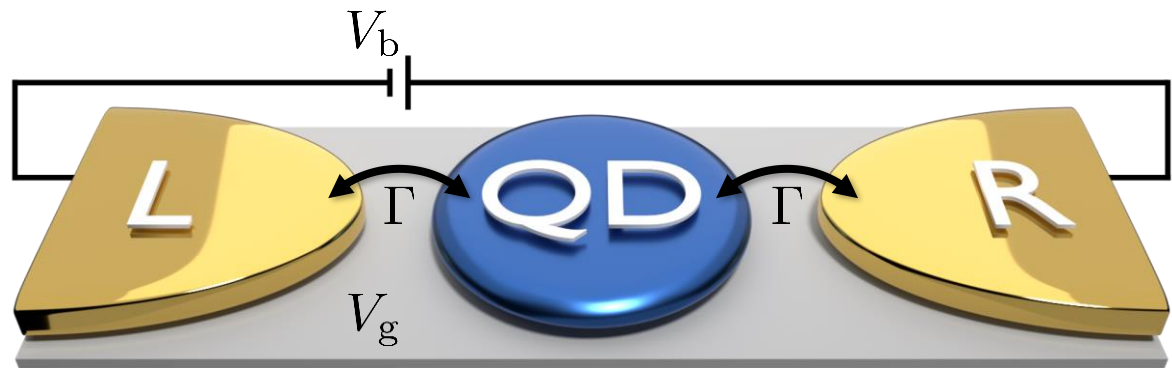
Statistics

Counting field:  $\chi$

$$\dot{\rho} = \mathcal{L}\rho$$

Full counting statistics

Current cumulants



Single electron transistor

Master equation:

$$0 = \mathcal{L}\rho^\infty = -\frac{i}{\hbar} \left[ \hat{H}_{\text{CNT}} + \hat{H}_{\text{LS}}, \rho^\infty \right] + \mathcal{L}_{\text{tun}}\rho^\infty + \mathcal{L}_{\text{rel}}\rho^\infty$$

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↑
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Reduced density matrix
   
 Lamb shift
   
 tunneling
   
 relaxation

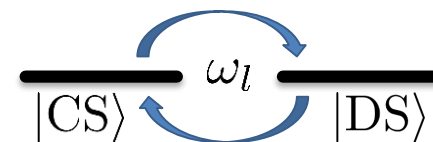
Minimal model for  $N = 0 \leftrightarrow 1$  (no extrinsic relaxation):

$$0 = \dot{\rho}_1 = -\frac{i}{\hbar} \left[ \hat{H}_{\text{LS}}, \rho_1 \right] + 2\Gamma_L \mathcal{R}_L \rho_0 - \frac{\Gamma_R}{2} \{ \mathcal{R}_R, \rho_1 \}$$

$$0 = \dot{\rho}_0 = \Gamma_R \text{tr} \{ \mathcal{R}_R \rho_1 \} - 4\Gamma_L \rho_0$$

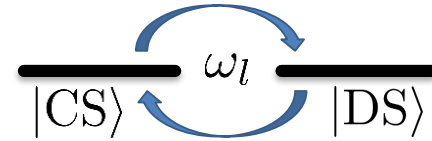
Lamb shift:

$$\hat{H}_{\text{LS}} = \frac{\hbar}{2} \sum_l \omega_l \mathcal{R}_l$$





$$\hat{H}_{\text{LS}} = \frac{\hbar}{2} \sum_l \omega_l \mathcal{R}_l$$



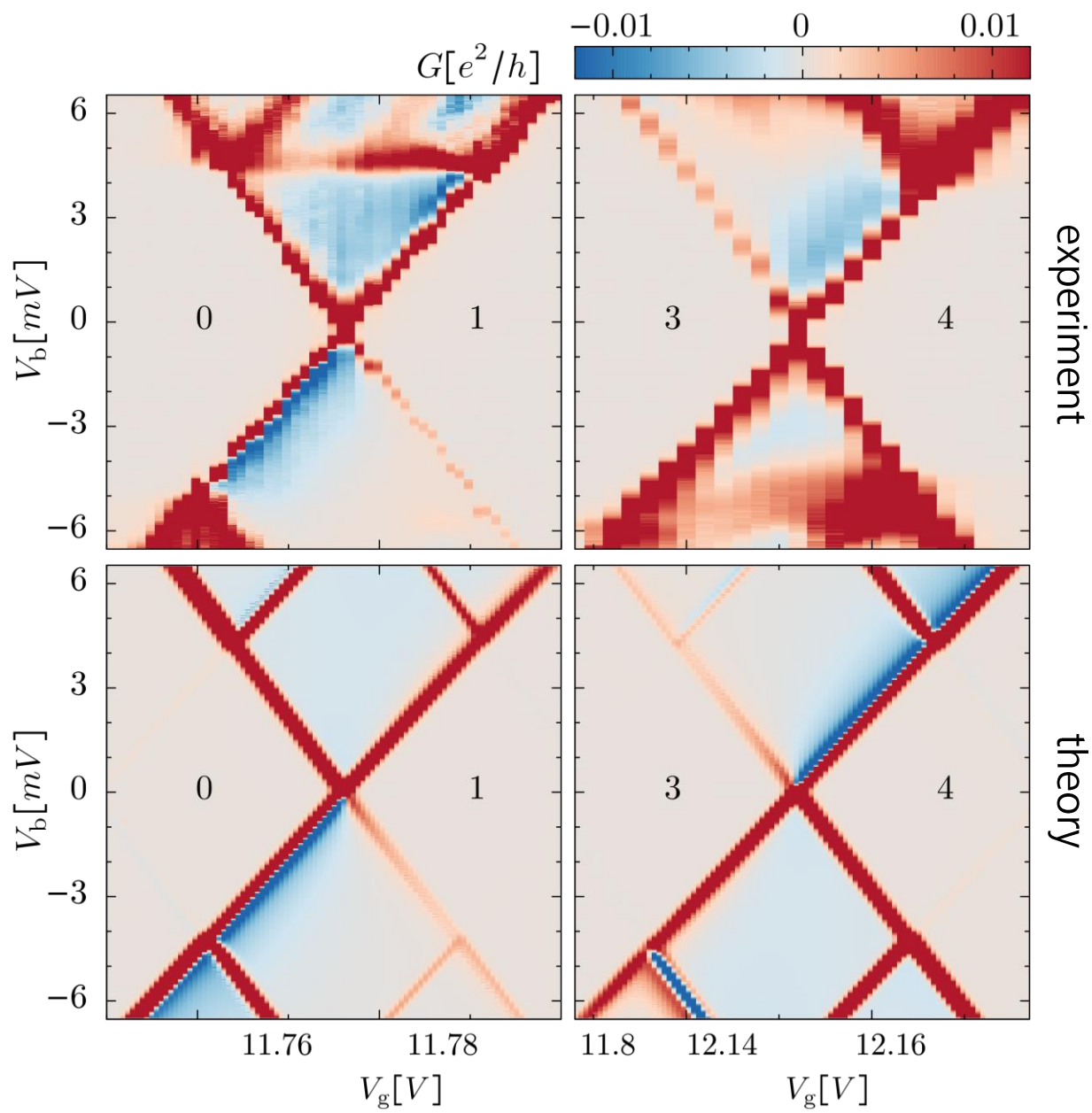
$$\omega_l(V_g, V_b) = \frac{\Gamma_l}{\pi} \left[ p_l(\alpha_g V_g) - p_l\left(U - \frac{J}{2} + \alpha_g V_g\right) \right]$$

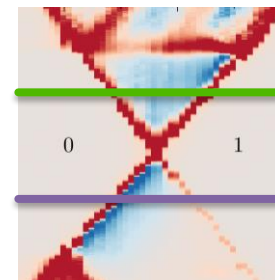
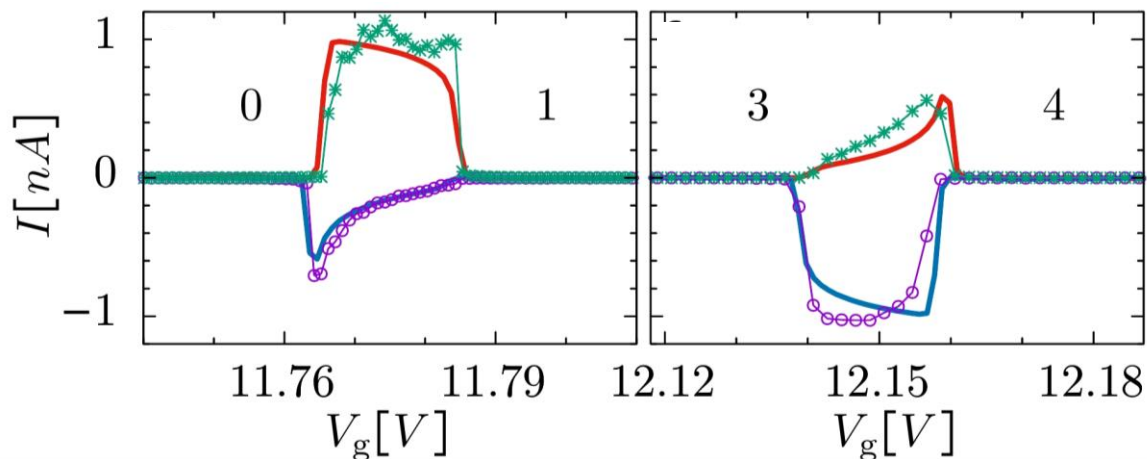
$$p_l(\Delta E) = -\text{Re} \psi \left( \frac{1}{2} + i \frac{\Delta E - \mu_l}{2\pi k_B T} \right) \quad \begin{array}{l} \mu_L = \eta e V_b \\ \mu_R = (\eta - 1) e V_b \end{array}$$

Digamma function

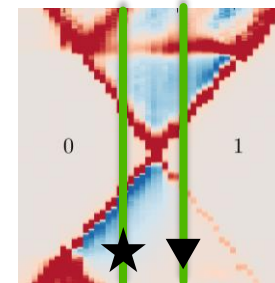
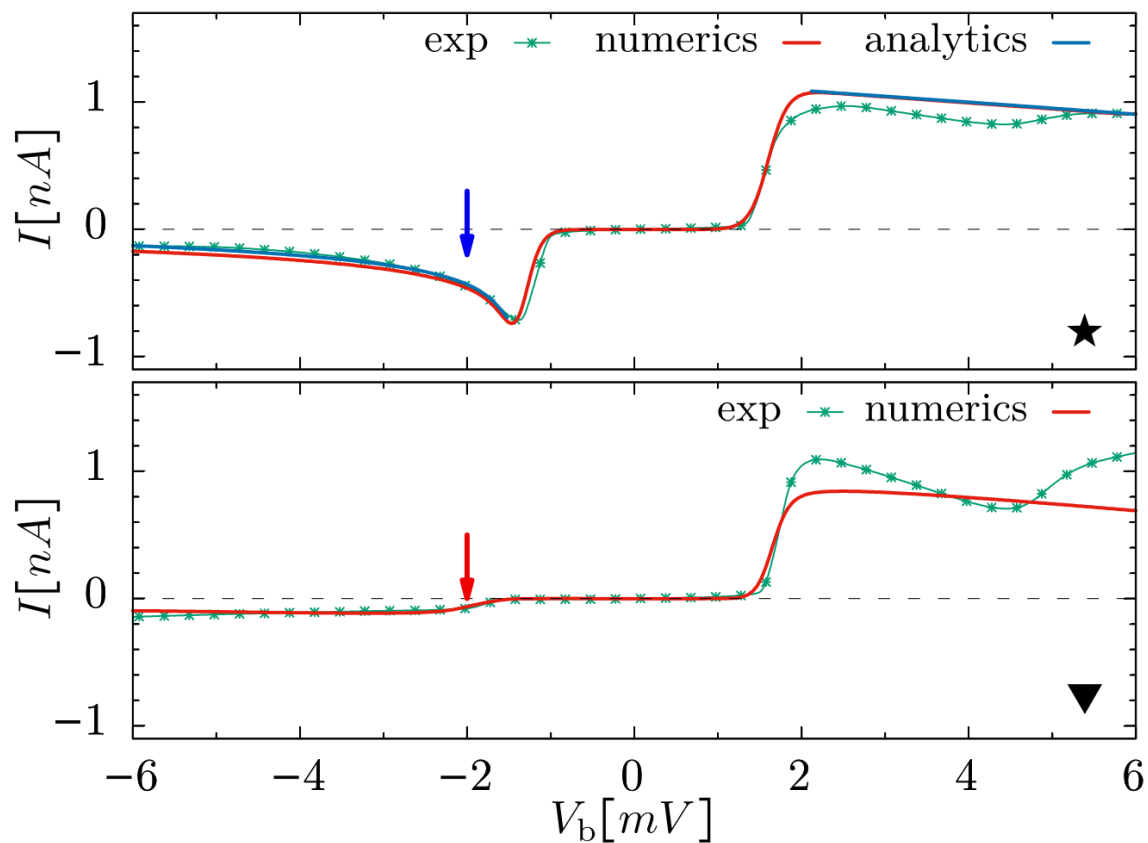
Terms come from principal value integrals over the Fermi functions:

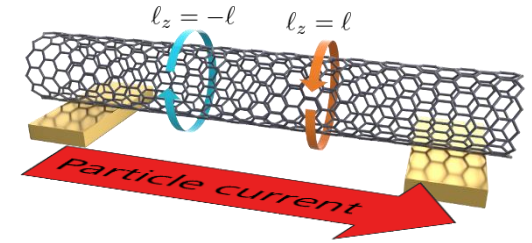
$$p.v. \int_{-\infty}^{\infty} d\epsilon \frac{f^{\pm}(\epsilon)}{\epsilon - \Delta E}$$





Bias traces:

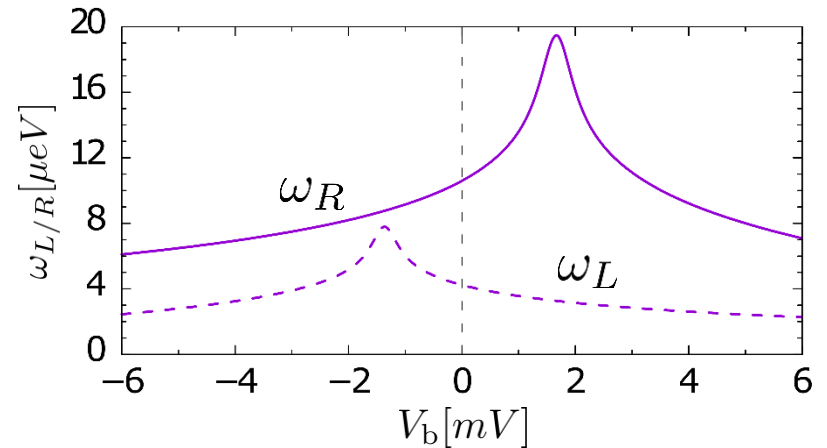
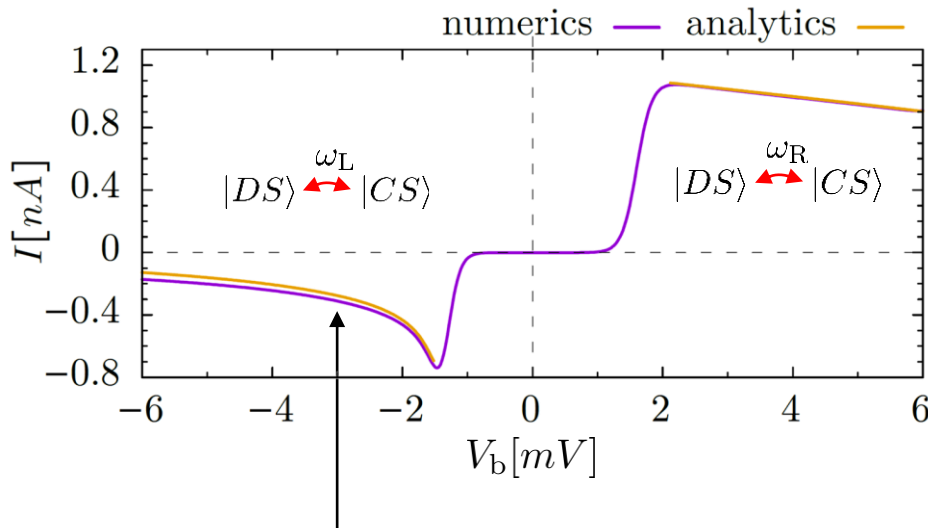




$$\dot{\rho} = -\frac{i}{\hbar}[H_{\text{CNT}}, \rho] - \frac{i}{\hbar}[H_{\text{LS}}, \rho] + \mathcal{L}_{\text{tun}}[\rho] + \mathcal{L}_{\text{rel}}[\rho]$$

$$H_{\text{LS}} = \frac{\hbar}{2}(\omega_L \mathcal{R}_L + \omega_R \mathcal{R}_R)$$

$$[\mathcal{R}_\alpha, |DS, \alpha\rangle\langle DS, \alpha|] = 0 \quad \rightarrow \quad \text{Only the **source** contribution perturbs the dark state}$$

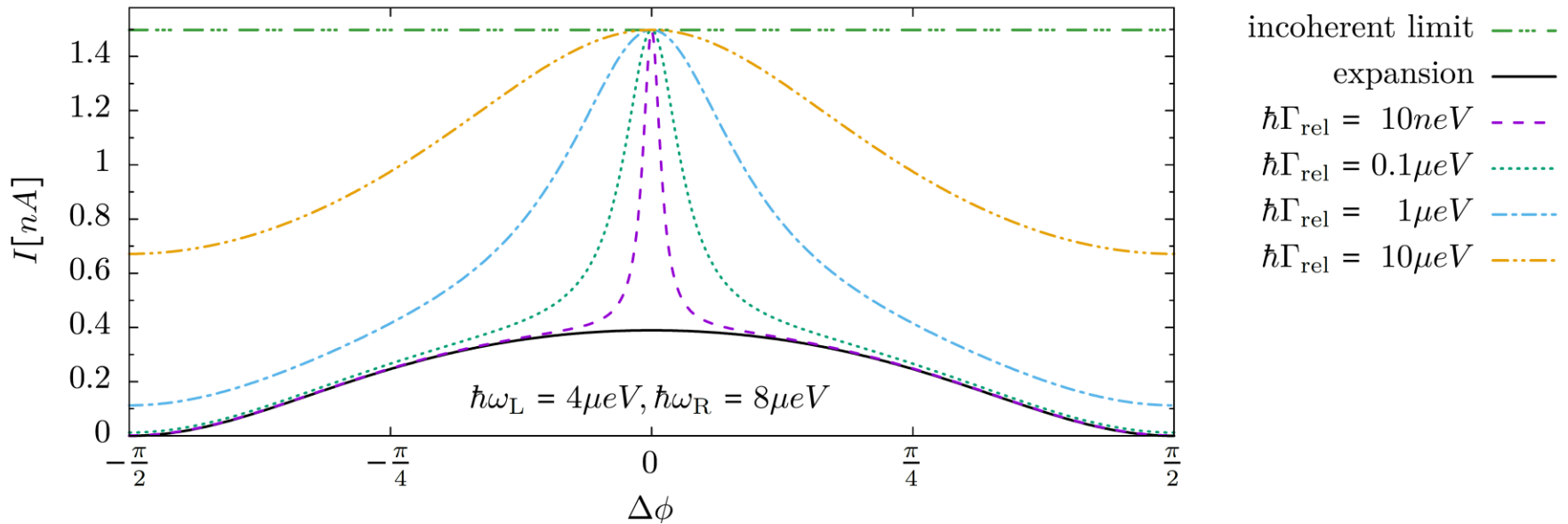


$$I = \frac{4e\Gamma_R\omega_L^2(1 + \cos \Delta\phi)}{\Gamma_R^2 + 4(\omega_L - \omega_R)^2 + \omega_L(1 + \cos \Delta\phi) [\omega_L\Gamma_R/\Gamma_L + 4\omega_R]}$$

$$\Delta\phi = \phi_L - \phi_R$$

$$I = \frac{4e\Gamma_R\omega_L^2(1 + \cos \Delta\phi)}{\Gamma_R^2 + 4(\omega_L - \omega_R)^2 + \omega_L(1 + \cos \Delta\phi) [\omega_L\Gamma_R/\Gamma_L + 4\omega_R]}$$

$$\Delta\phi = \phi_L - \phi_R$$



$$I(\Delta\phi = 0) = \frac{4e\Gamma_L\Gamma_R}{4\Gamma_L + \Gamma_R}$$

incoherent current through 4 degenerate levels

$$H_{\text{CNT}} = \sum_{ml_z} (m\epsilon_0 - \xi) \hat{n}_{ml_z} + \frac{U}{2} \hat{N}^2 + J \sum_m \left( \hat{\mathbf{S}}_{ml} \cdot \hat{\mathbf{S}}_{m-l} + \frac{1}{4} \hat{n}_{ml} \hat{n}_{m-l} \right)$$

$$J = 0$$

2-particle groundstate is 6-fold degenerate  
 → new dark state possible

$$|2, \text{DS}\rangle = \frac{1}{2} \left( e^{2il\phi_\alpha} \begin{array}{cc} \text{---} & \text{---} \\ \uparrow\downarrow & \text{---} \\ \uparrow\downarrow & \uparrow\downarrow \end{array} - \begin{array}{cc} \text{---} & \text{---} \\ \uparrow & \downarrow \\ \uparrow\downarrow & \uparrow\downarrow \end{array} + \begin{array}{cc} \text{---} & \text{---} \\ \downarrow & \uparrow \\ \uparrow\downarrow & \uparrow\downarrow \end{array} + e^{-2il\phi_\alpha} \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \uparrow\downarrow \\ \uparrow\downarrow & \uparrow\downarrow \end{array} \right)$$

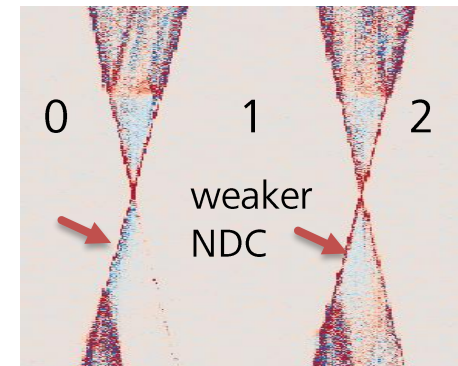
→ interference NDC also at  $1 \leftrightarrow 2$  particle resonance

$$J > \Gamma$$

2-particle groundstate is the inter-valley singlet

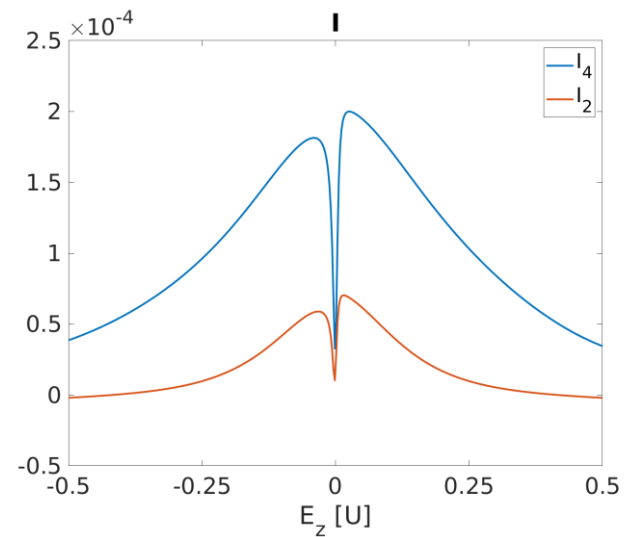
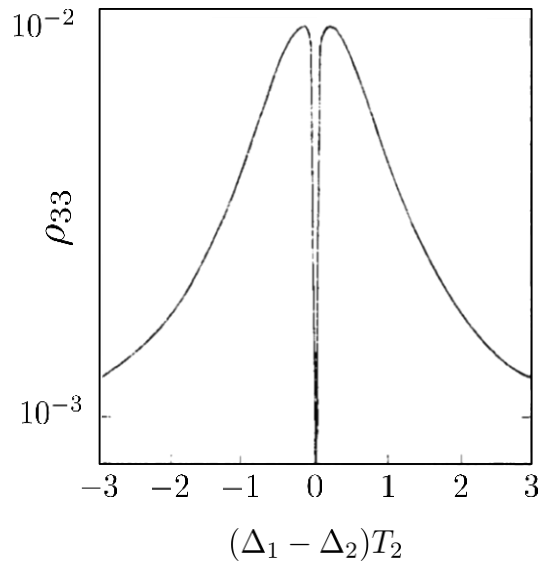
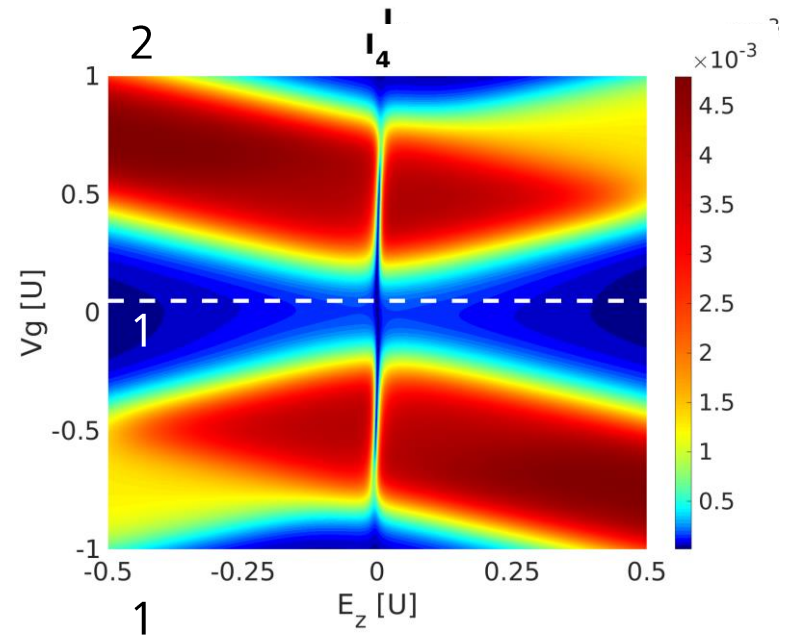
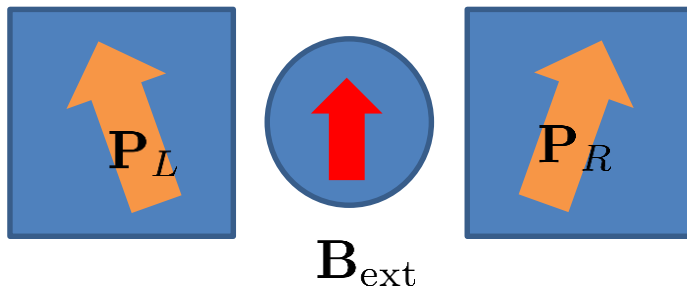
$$|2_0\rangle = \frac{1}{\sqrt{2}} \left( \begin{array}{cc} \text{---} & \text{---} \\ \uparrow & \downarrow \\ \uparrow\downarrow & \uparrow\downarrow \end{array} - \begin{array}{cc} \text{---} & \text{---} \\ \downarrow & \uparrow \\ \uparrow\downarrow & \uparrow\downarrow \end{array} \right)$$

→ no 2 particle dark state is possible



The general tunnelling matrix for a degenerate two level system

$$\Gamma^\alpha = \Gamma^\alpha (\mathbb{1} + \mathbf{P}_\alpha \cdot \boldsymbol{\sigma})$$





# Acknowledgments

M. Niklas



M. Grifoni



C. Rohrmeier



and thanks to ...



M. Schafberger



N. Paradiso

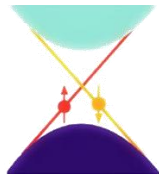


C. Strunk

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SFB 689



IDK Top. Iso.

Elitenetzwerk  
Bayern



GRK 1570



**SFB 1277**  
Collaborative Research Center  
Universität Regensburg





Interference occurs between the **degenerate** angular momentum states of a **zig-zag class** carbon nanotube

Spatially confined tunneling mixes angular momentum states providing a non-diagonal **tunnelling rate matrix**

The **dark state** is the eigenvector of the tunneling rate matrix with zero eigenvalue

The **Lamb-shift** precession perturbs the dark state and explains the **bias voltage asymmetry**

**Thank you for your attention**