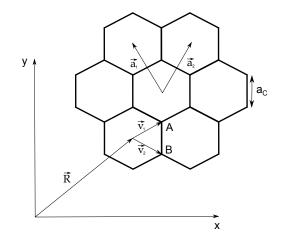
## Quantum Theory of Condensed Matter I

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5.1.01 Mondays 10:15 9.2.01 Tuesdays 12:15

Sheet 4

## 1. Electronic structure of non-interacting graphene



A graphene sheet is a honeycomb lattice of carbon atoms, as shown above. Let the distance between carbon atoms be  $a_C = 1.42$  Å. A good model for graphene is to consider a plane in which there is one valence electron per carbon atom. To calculate the electronic bands of the non-interacting problem we will use the LCAO method with a single  $p_z$  orbital at each carbon site. Let  $\vec{R}$  denote the centers of the hexagons in the honeycomb: these form the underlying Bravais lattice. (Note that it is different from the graphene hexagonal lattice.) The unit cell spanned by  $\vec{a}_1$  and  $\vec{a}_2$  contains two carbon atoms conventionally labelled as A and B atoms, located at  $\vec{R} + \vec{\tau}_A$ ,  $\vec{R} + \vec{\tau}_B$  as shown in the figure. The tight-binding amplitude connecting these sites is denoted by -t. We can choose the zero of energy so that the energy of an isolated  $p_z$  orbital is zero. Then the entire Hamiltonian consists only of the nearest-neighbour hopping.

- 1. Find the reciprocal lattice and construct the first Brillouin zone. (1 Point)
- 2. Write the Hamiltonian in the basis of the extended states

$$\begin{split} \left| \phi_{A\vec{k}} \right\rangle &= \frac{1}{\sqrt{N_{cell}}} \sum_{\vec{R}} e^{i\vec{k}\cdot\vec{R}} \left| p_z, \vec{R} + \vec{\tau}_A \right\rangle \\ \left| \phi_{B\vec{k}} \right\rangle &= \frac{1}{\sqrt{N_{cell}}} \sum_{\vec{R}} e^{i\vec{k}\cdot\vec{R}} \left| p_z, \vec{R} + \vec{\tau}_B \right\rangle, \end{split}$$

where  $N_{cell}$  is the number of the unit cells of the crystal and  $|p_z, \vec{S}\rangle$  is the state of the  $p_z$  orbital localized in the position  $\vec{S}$ . Hint: make use of the relation

$$\frac{1}{N_{cell}}\sum_{\vec{R}}\exp\left\{i(\vec{k}-\vec{k}')\cdot\vec{R}\right\} = \delta_{\vec{k}\vec{k}'},$$

where the sum is intended over the Bravais lattice and  $\vec{k}$  and  $\vec{k'}$  belong to the reciprocal lattice.

(2 Points)

3. The generic extended states of the graphene sheet have the form:

 $|\psi\rangle = \sum_{\vec{k}} c_{A\vec{k}} |\phi_{A\vec{k}}\rangle + c_{B\vec{k}} |\phi_{B\vec{k}}\rangle$ , where  $c_{A\vec{k}}$  and  $c_{B\vec{k}}$  are complex numbers. Find the eigenstates of the system and the corresponding eigenvalues. How many bands do you find? What are the energy-wavevector relations? (2 Points)

Hint: You will need to solve a  $2 \times 2$  matrix eigenvalue problem - the time-independent Schrödinger equation written in the basis of extended states.

- 4. Show that the Fermi energy is equal to zero by verifying that this gives the correct electron density. Find the set of  $\vec{k}$  points for which  $\epsilon(\vec{k}) = 0$ . Show that these correspond to the corners of the first Brillouin zone. (2 Points)
- 5. Plot the complete dispersion relation of graphene.

6. (Oral) Show that near to a (first) Brillouin zone corner with wavevector  $\vec{K}$  (note that  $\vec{K}$  is the location of a BZ corner, not a reciprocal lattice vector), the spectrum is approximately

$$\epsilon(\vec{k}) \approx \pm \hbar v \sqrt{(k_x - K_x)^2 + (k_y - K_y)^2}$$

Determine the velocity v in terms of t and  $a_C$ . This behaviour is intermediate between that of a metal (with a Fermi surface instead of Fermi points  $\vec{K}$ ) and an insulator (with a band gap). Indeed, graphene behaves like a semi-metal, with poor conductivity yet significantly higher than in an insulator.

## **Frohes Schaffen!**