

Mesoscopic Physics

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Wednesdays at 15:15

Sheet 2

1. Naïve distribution function

Consider the trial distribution function given by $f(q, p) = \text{Tr} \{ \hat{\rho} \delta(q - \hat{q}) \delta(p - \hat{p}) \}$. Prove that this function is complex and thus represents a bad choice.

Hint: It is enough to provide a counter example. Consider for instance the pure state density operator defined by $\hat{\rho} = |\psi\rangle\langle\psi|$, built from a wave function such as $\psi(x) = \sqrt{\alpha/\pi} \sin(\frac{\alpha x}{\hbar})$, with x defined in the interval $[0, 2\pi\hbar/\alpha]$.

2. Wigner distribution

- The Wigner distribution function is defined as

$$f_W(q, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dy \left\langle q - \frac{y}{2} \left| \hat{\rho} \right| q + \frac{y}{2} \right\rangle e^{i\frac{py}{\hbar}}$$

Prove that it has the following properties:

(i) It is real. That is, $\text{Im}(f_W) = \frac{1}{2}(f_W - f_W^*) = 0$.

(ii) f_W satisfies:

$$\begin{aligned} \int dp f_W(q, p) &= \langle q | \hat{\rho} | q \rangle \\ \int dq f_W(q, p) &= \langle p | \hat{\rho} | p \rangle \\ \int dq dp f_W(q, p) &= \text{Tr}(\hat{\rho}) = 1 \end{aligned}$$

Remember that $2\pi\delta(x) = \int_{-\infty}^{\infty} e^{i\kappa x} d\kappa$ and that $\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ixp/\hbar}$.

(iii) f_W is Galilei invariant:

$$\begin{aligned} \psi(q) \rightarrow \psi(q+a) &\implies f_W(q, p) \rightarrow f_W(q+a, p) \\ \psi(q) \rightarrow e^{ip'q/\hbar} \psi(q) &\implies f_W(q, p) \rightarrow f_W(q, p-p') \end{aligned}$$

Hint: For the density operator $\hat{\rho} = \sum_n p_n |\psi_n\rangle\langle\psi_n|$, we can rewrite the Wigner distribution function as

$$f_W(q, p) = \frac{1}{2\pi\hbar} \sum_n p_n \int_{-\infty}^{\infty} dy \psi_n\left(q - \frac{y}{2}\right) \psi_n^*\left(q + \frac{y}{2}\right) e^{i\frac{py}{\hbar}}.$$

Calculate $\overline{f_W}$ associated to $\overline{\psi_n}(q) = \psi_n(q+a)$ and $\overline{\psi_n}(q) = e^{i\frac{p'q}{\hbar}} \psi_n(q)$ respectively.

(iv) f_W is invariant under space and time reflections:

$$\begin{aligned}\psi(q) \rightarrow \psi(-q) &\implies f_W(q, p) \rightarrow f_W(-q, -p) \\ \psi(q) \rightarrow \psi^*(q) &\implies f_W(q, p) \rightarrow f_W(q, -p)\end{aligned}$$

(v) In the force-free case the equation of motion is the classical one:

$$\frac{\partial f_W}{\partial t} = -\frac{p}{m} \frac{\partial f_W}{\partial q}$$

Consider that the time derivative of the density operator is given by $\dot{\hat{\rho}} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}]$ and if no forces are present we can write the Hamiltonian as $\hat{H} = \hat{p}^2/2m$.

3. Weyl representation

Prove that the Weyl representation for $F = q^n p^m$ is

$$\hat{F}_{Weyl} = \frac{1}{2^n} \sum_{r=0}^n \binom{n}{r} \hat{q}^{n-r} \hat{p}^m \hat{q}^r$$

Hint: The associated operator for the classical quantity F is given by Weyl's expansion:

$$\hat{F}_{Weyl}(\hat{q}, \hat{p}) = \int d\sigma d\tau e^{\frac{i}{\hbar}(\sigma\hat{q} + \tau\hat{p})} \varphi(\sigma, \tau)$$

where

$$\varphi(\sigma, \tau) = \frac{1}{(2\pi\hbar)^2} \int dq dp e^{\frac{i}{\hbar}(\sigma q + \tau p)} q^n p^m$$

Frohes Schaffen!