

Density Matrix Theory

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5.0.21, Tuesdays, 10:15

Sheet 6

1. Equilibrium: the free energy formulation

Consider the master equation for the Anderson impurity model introduced in the Sheet 4:

$$\begin{aligned}\dot{P}_0 &= -2\gamma f^+(\varepsilon_d)P_0 + \gamma \sum_{\sigma} f^-(\varepsilon_d)P_{1\sigma} \\ \dot{P}_{1\sigma} &= -\gamma[f^+(\varepsilon_d + U) + f^-(\varepsilon_d)]P_{1\sigma} \\ &\quad + \gamma f^+(\varepsilon_d)P_0 + \gamma f^-(\varepsilon_d + U)P_2 \\ \dot{P}_2 &= -2\gamma f^-(\varepsilon_d + U)P_2 + \gamma \sum_{\sigma} f^+(\varepsilon_d + U)P_{1\sigma}\end{aligned}$$

where

$$P_0(t) \equiv \langle 0 | \rho_{red}(t) | 0 \rangle, P_{1\sigma} \equiv \langle 1\sigma | \rho_{red}(t) | 1\sigma \rangle, P_2(t) \equiv \langle 2 | \rho_{red}(t) | 2 \rangle$$

are the populations of the reduced density matrix with respect to the manybody energy eigenbasis $|0\rangle, |1\uparrow\rangle, |1\downarrow\rangle, |2\rangle$ of the impurity.

1. Prove that the stationary solution of this master equation is independent of the magnitude of the bare tunnelling rate γ and, for every value of the parameters $(\varepsilon_d, U, \mu, T)$ defining the model, can be written in the form:

$$\begin{aligned}P_0^{stat} &= \frac{1}{N} f^-(\varepsilon_d) f^-(\varepsilon_d + U) \\ P_{1\sigma}^{stat} &= \frac{1}{N} f^+(\varepsilon_d) f^-(\varepsilon_d + U) \\ P_2^{stat} &= \frac{1}{N} f^+(\varepsilon_d) f^+(\varepsilon_d + U)\end{aligned}\tag{1}$$

where N is the normalization factor that ensures the sum of the probability to be 1. Moreover $f^+(\varepsilon) \equiv [1 + e^{\beta(\varepsilon - \mu)}]^{-1}$ and $f^-(\varepsilon) \equiv 1 - f^+(\varepsilon)$.

2. Prove that the equilibrium probabilities derived at the previous point can be obtained from a thermodynamical formulation of the problem where the impurity, defined by the Hamiltonian H_S (see Sheet 4), can exchange energy and particles with a bath with temperature T and chemical potential μ . In particular calculate the grand canonical partition function $\mathcal{Z} = \text{Tr}_S\{e^{-\beta(H_S - \mu N_S)}\}$ for the impurity and prove that:

$$P_{\alpha}^{stat} = \frac{1}{\mathcal{Z}} \text{Tr}_S\{|\alpha\rangle\langle\alpha| e^{-\beta(H_S - \mu N_S)}\}$$

where $|\alpha\rangle$ is a manybody energy eigenstate of the impurity and N_S the particle number.

2. Time evolution for a Markovian master equation

In this exercise we consider the Markoff master equation (1) and calculate numerically the time evolution for the populations of the many-body states of the impurity.

1. Show that the equations (1) can be cast into a matrix form $\dot{P}(t) = LP(t)$ where $P \equiv (P_0, P_{1\uparrow}, P_{1\downarrow}, P_2)^T$ and

$$L = \gamma \begin{pmatrix} -2f^+(\varepsilon_d) & f^-(\varepsilon_d) & f^-(\varepsilon_d) & 0 \\ f^+(\varepsilon_d) & -f^-(\varepsilon_d) - f^+(\varepsilon_d + U) & 0 & f^-(\varepsilon_d + U) \\ f^+(\varepsilon_d) & 0 & -f^-(\varepsilon_d) - f^+(\varepsilon_d + U) & f^-(\varepsilon_d + U) \\ 0 & f^+(\varepsilon_d + U) & f^+(\varepsilon_d + U) & -2f^-(\varepsilon_d + U) \end{pmatrix}.$$

Prove that the solution of the equation can be written in the form $P(t) = e^{Lt}P(t=0)$. Taking advantage of this algebraic formulation, calculate the numerical solution of (1). Hint: the function “expm” calculates the exponential of a matrix in Matlab.

2. Prove that, if the time is measured in units of $1/\gamma$ solutions with different tunneling rates coincide.
3. Check that the stationary solution is reached by the system after a time corresponding to a few $1/\gamma$ and that it is independent of the initial condition.
4. Calculate the time evolution for the population vector P also with the help of one of the packages for ordinary differential equations available in Matlab. Compare the results with the previous method. Hint: There are different types of solvers. You can start by typing “help ode23” in the command line and read the documentation.

Frohes Schaffen!